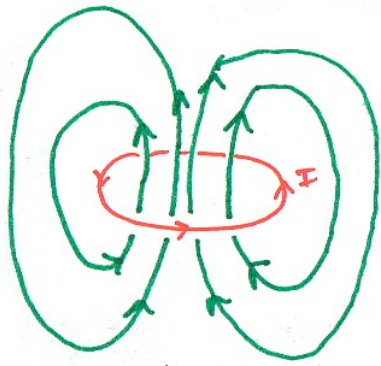


Model for permanent magnets



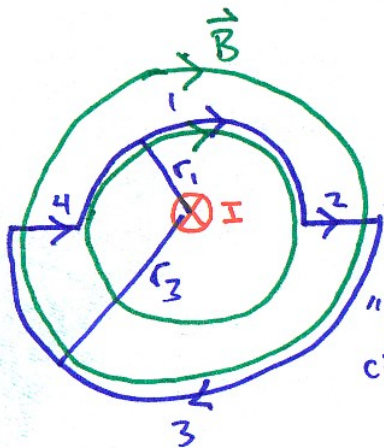
The field due to a current loop looks the same as that of a permanent magnet. It is now well established that

permanent magnetism is due to the tiny "current loops" associated with electron spin.

Ampère's law

Book: Biot-Savart

$$\rightarrow B_{\text{wire}} = \frac{\mu_0 I}{2\pi r}$$



Integration path

$$\oint \vec{B} \cdot d\vec{l} = \int_1 + \int_2 + \int_3 + \int_4$$

"magnetic circulation"

= 0 because $\vec{B} \perp d\vec{l}$

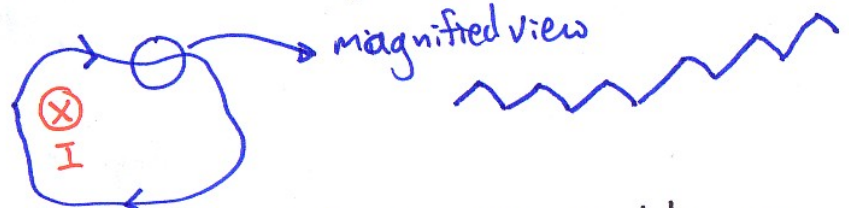
For 1 & 3, $\vec{B} \parallel d\vec{l} \Rightarrow \oint \vec{B} \cdot d\vec{l} = \int_1 B dl + \int_3 B dl$

B is constant along 1 & 3

$$\Rightarrow \oint \vec{B} \cdot d\vec{l} = \frac{\mu_0 I}{2\pi r_1} \int_1 dl + \frac{\mu_0 I}{2\pi r_3} \int_3 dl$$

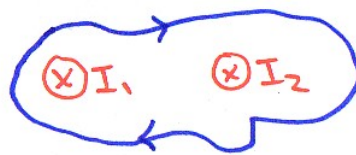
= $\mu_0 I$ Note: no r-dependence!

We can make any loop by combining radial & circumferential pieces:



$$\Rightarrow \oint \vec{B} \cdot d\vec{l} = \mu_0 I \text{ for any loop around the wire}$$

What if there are two wires?



$$\vec{B}_{\text{TOT}} = \vec{B}_1 + \vec{B}_2$$

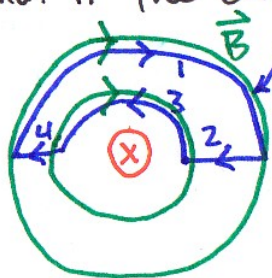
\vec{B} due to I_1 \vec{B} due to I_2

$$\oint \vec{B}_{\text{TOT}} \cdot d\vec{l} = \oint (\vec{B}_1 + \vec{B}_2) \cdot d\vec{l} = \oint \vec{B}_1 \cdot d\vec{l} + \oint \vec{B}_2 \cdot d\vec{l}$$

$$= \mu_0 I_1 + \mu_0 I_2$$

$$\Rightarrow \oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{net}}$$

What if the current doesn't thread the loop?



Integration path

Along 1, $\vec{B} \parallel d\vec{l} \Rightarrow \int_1 \vec{B} \cdot d\vec{l} = \int_1 B dl = \frac{\mu_0 I}{2}$

Along 3, $\vec{B} \nabla d\vec{l}$
anti-parallel $\Rightarrow \int_3 \vec{B} \cdot d\vec{l} = -\int_3 B dl = -\frac{\mu_0 I}{2}$

$\Rightarrow \oint \vec{B} \cdot d\vec{l} = 0 \Rightarrow I$ only "counts" toward $\oint \vec{B} \cdot d\vec{l}$ if it threads the loop.

Putting this all together:

$$\oint \vec{B} \cdot d\vec{l} = \mu_0 I_{\text{net}} + \text{term that depends on changing } \vec{E} \text{ threading}$$

Ampère's Law: The third MEQ